

Describing Pseudospherical Planes and Other Properties of Evolutionary Soliton Equations

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Abstract: In this paper we will derive Bäcklund transformations and conservation laws based on geometrical properties of evolution equations with more than two independent variables that describe pseudospherical surfaces.

Keywords: Evolution equations, Pseudospherical surfaces, Bäcklund transformations, conservation laws and Solitons.

I. Introduction

In this paper we, interest in Bäcklund transformation [12], and its connection with some special equations and their associated soliton theory. Under this transformation an infinite family of constant curvature surfaces can be produced from a given one. The notion of a differential equation for a function $u(x, t)$ that describes a pseudospherical surface (P.S.S.) was introduced in [1,6,7], where classifications for some equations of types

$$u_{xt} = \psi \left(u, u_x, u_{xx}, \dots, \frac{\partial^k u}{\partial x^k} \right) \quad \text{and} \quad u_t = \psi \left(u, u_x, \dots, \frac{\partial^k u}{\partial x^k} \right)$$

Were obtained. Furthermore characterizations of equations with more than two independent variables of types

$$u_{xt} = \psi \left(u, u_x, \dots, \frac{\partial^k u}{\partial x^k}, u_y, \dots, \frac{\partial^k u}{\partial x^k} \right), \quad u_t = \psi \left(u, u_x, \dots, \frac{\partial^k u}{\partial x^k}, u_y, \dots, \frac{\partial^k u}{\partial x^k} \right)$$

$$\text{and } u_{tt} = \psi \left(u, u_x, \dots, \frac{\partial^k u}{\partial x^k}, u_y, \dots, \frac{\partial^k u}{\partial y^k}, u_t \right) \text{ are given in [2,3,4].}$$

A systematic procedure to determine linear problems associated to non-linear equations of the above types was also introduced in case of two independent variables.

In this work, we consider evolution equations for a function $u(x, y, t)$ that describes an (η, ξ) 3-dim. P.S.P. as given in [2,3,4] and we investigate an analogous method to derive Bäcklund transformations and conservation laws based on geometrical properties of these 3-dimensional pseudo spherical planes in R^5 .

II. Local theory of constant negative curvature submanifolds of R^{2n-1}

Let M be an n -dimensional Riemannian manifold with constant curvature K isometrically immersed in \bar{M}^{2n-1} with constant curvature \bar{K} , with $K < \bar{K}$. Let $e_1, e_2, \dots, e_{2n-1}$ be a moving orthonormal frame on an open set of \bar{M} , so that at points of M, e_1, e_2, \dots, e_n are tangents to M . Let ω_A be the dual orthonormal coframe and consider ω_{AB} defined by [2]

$$de_A = \sum_B \omega_{AB} e_B$$

The structure equations of \bar{M} are

$$d\omega_A = \sum_B \omega_B \wedge \omega_{BA}, \quad \omega_{AB} + \omega_{BA} = 0 \quad (1)$$

$$d\omega_{AB} = \sum_C \omega_{AC} \wedge \omega_{CB} - \bar{K} \omega_A \wedge \omega_B \quad \text{with} \quad 1 \leq A, B, C \leq 2n-1 \quad (2)$$

Restricting these forms to M we have $\omega_\alpha = 0$, so (1) gives with $n+1 \leq \alpha, \beta, \gamma \leq 2n-1$ and $1 \leq I, J, L \leq n$,

$$d\omega_\alpha = \sum_I \omega_I \wedge \omega_{I\alpha} = 0 \quad (3)$$

$$d\omega_I = \sum_J \omega_J \wedge \omega_{JI} \quad (4)$$

from (2) we obtain, Gauss equation

$$d\omega_{ij} = \sum_L \omega_{iL} \wedge \omega_{Lj} + \sum_\alpha \omega_{i\alpha} \wedge \omega_{\alpha j} - \bar{K} \omega_i \wedge \omega_j \quad (5)$$

and Codazzi equation

$$d\omega_{i\alpha} = \sum_A \omega_{iA} \wedge \omega_{A\alpha} \quad (6)$$

M has constant sectional curvature K if and only if

$$\Omega_{ij} = d\omega_{ij} - \sum_L \omega_{iL} \wedge \omega_{Lj} = -K \omega_i \wedge \omega_j \quad (7)$$

$$\sum_\alpha \omega_{i\alpha} \wedge \omega_{\alpha j} = (\bar{K} - K) \omega_i \wedge \omega_j \quad (8)$$

Also, equation (2) implies that [2]

$$d\omega_{\alpha\beta} = \sum_\gamma \omega_{\alpha\gamma} \wedge \omega_{\gamma\beta} + \Omega_{\alpha\beta}$$

With $\Omega_{\alpha\beta} = \sum_I \omega_{\alpha I} \wedge \omega_{I\beta} \quad (9)$

The forms $\Omega_{\alpha\beta}$ give the normal curvature of M and $I = \sum_I (\omega_I)^2$ is its first fundamental form.

For our purpose in this paper, we write these equations when \bar{M} is taken to be R^5 and M is a 3-dimensional submanifold with constant sectional curvature $K = -1$ (i.e. pseudo spherical 3-plane in R^5).

The equations take the forms [2]

$$\left. \begin{aligned} d\omega_1 &= \omega_4 \wedge \omega_2 + \omega_5 \wedge \omega_3 \\ d\omega_2 &= -\omega_4 \wedge \omega_1 + \omega_6 \wedge \omega_3 \\ d\omega_3 &= -\omega_5 \wedge \omega_1 - \omega_6 \wedge \omega_2 \\ d\omega_4 &= \omega_1 \wedge \omega_2 \\ d\omega_5 &= \omega_1 \wedge \omega_3 \\ d\omega_6 &= \omega_2 \wedge \omega_3 \end{aligned} \right\} \quad (10)$$

where we have written

$$\begin{aligned} \omega_4 &= \omega_{12} & \omega_5 &= \omega_{13}, \text{ and} \\ \omega_6 &= \omega_{23} \text{ with } & \omega_{ij} &= -\omega_{ji}, i, j = 1, 2, 3, \quad \omega_{ii} = 0 \end{aligned}$$

We shall recall here the definition of a differential equation to describe a pseudospherical surface, introduced in [1] and modify it in order to suit our purposes here.

Definition 2.1

A differential equation E -for a real function $u(x, y, t)$ describes a 3-dimensional pseudospherical plane in R^5 (simply P.S.P.) if it is the necessary and sufficient condition for the existence of differentiable functions $f_{\alpha i}, 1 \leq \alpha \leq 6$ and $1 \leq i \leq 3$, depending on u and its derivatives, such that the 1-forms [2,3]

$$\omega_\alpha = f_{\alpha 1} dx + f_{\alpha 2} dy + f_{\alpha 3} dt \quad (11)$$

satisfy the structure equations of a 3-plane of constant sectional curvature -1 in R^5 i.e. equations (10).

Definition 2.2

We shall define such 3-dimensional P.S.P to be a two-parameters 3-dimensional P.S.P. $f_{31} = f_{41} = \eta$ and $f_{22} = f_{42} = \xi$, with η and ξ constant parameters. In Fact, one can see that when $u(x, y, t)$ is a generic solution of E , it provides a metric defined on an open subset of R^3 , whose sectional curvature is -1 and the lengths of the vector fields $\frac{\partial}{\partial x}$ and $\frac{\partial}{\partial y}$ satisfy $\left| \frac{\partial}{\partial x} \right|^2 \geq \eta^2, \left| \frac{\partial}{\partial y} \right|^2 \geq \xi^2$. [2,3]

III. Generalization of Bäcklund's theorem

In this section, we define a pseudospherical geodesic congruence between two n -dimensional submanifolds M and M' of a space form \bar{M}_k^{2n-1} with constant sectional curvature K . We prove a generalization of Bäcklund's theorem, [12] for such submanifolds and the complete integrability of the differential ideal associated to the existence of a pseudospherical congruence.

In what follows we need the notion of angles between two k -planes in a $2k$ -dimensional inner product space. [10]

Definition 3.1

Let E_1 and E_2 be two k -planes in a $2k$ -dimensional inner product space (V, \langle, \rangle) and $\pi: V \rightarrow E_1$ the orthogonal projection. Define a symmetric bilinear form on E_2 by $(v_1, v_2) = \langle P(v_1), P(v_2) \rangle$. The k angles between E_1 and E_2 are defined to be $\theta_1, \dots, \theta_k$ where $\cos^2 \theta_1, \dots, \cos^2 \theta_k$ are the k -eigenvalues for the self-adjoint operator $A: E_1 \rightarrow E_2$ such that $(v_1, v_2) = \langle Av_1, v_2 \rangle$. [10]

Definition 3.2

Suppose the n angles between E_1 and E_2 are $\theta_1, \dots, \theta_n$. Then it follows from the definition that there are two orthonormal bases e_1, \dots, e_{2n} and e'_1, \dots, e'_{2n} of V such that e'_1, \dots, e'_n are eigenvectors of A with eigenvalues $\cos^2 \theta_1, \dots, \cos^2 \theta_n$ respectively, e_1, \dots, e_n form a base for E_1 , and

$$\begin{aligned} e_i &= \cos \theta_i e_i + \sin \theta_i e_{n+i-1}, \\ e'_{n+i-1} &= -\sin \theta_i e_i + \cos \theta_i e_{n+i-1} \end{aligned}$$

for $1 \leq i \leq n$. [9]

Definition 3.3

A geodesic congruence between two n -dimensional submanifolds M and M' of a $(2n-1)$ -dimensional space form \bar{M} is a diffeomorphism $\ell: M \rightarrow M'$, such that for $P \in M$ and $P' = \ell(P)$, there exists a unique geodesic γ in \bar{M} joining P and P' , whose tangent vectors at P and P' are in $T_P M$ and $T_{P'} M'$ respectively. [10]

Definition 3.4

A geodesic congruence $\ell: M \rightarrow M'$ between two n -dimensional submanifolds of \bar{M} is called pseudospherical if:

- (1) the distance between P and $P' = \ell(P)$ on \bar{M} , is a constant r , independent of P ;
- (2) the $(n-1)$ angles between v_P and $v_{P'}$ are all equal to a constant θ , independent, of P ;
- (3) the normal bundles v and v' are flat ;
- (4) the bundle map $\Gamma: v \rightarrow v'$ given by the orthogonal projection commutes with the normal connections. [10]

Definition 3.5

For given geodesic congruence $\ell: M \rightarrow M'$, we remark that, the normal spaces v_P and $v_{P'}$, at corresponding points P and P' are $(n-1)$ dimensional and orthogonal to the plane determined by the position vector X of M and the tangent vector of γ at P . Therefore, v_P and $v_{P'}$, lie in a $(2n-2)$ dimensional vector space, i.e. there are $(n-1)$ angles between v_P and $v_{P'}$. [10]

Theorem 3.1

Suppose there is a pseudo-spherical congruence $l: M \rightarrow M'$ of n -manifolds in R^{2n-1} with distance r between corresponding points and angle $\theta \neq 0$ between corresponding normals. Then both M and M' have constant sectional curvature $-\sin^2 \theta / r^2$. [9]

Proof.

Since v' is flat, we may choose an orthonormal frame $e'_{n+1}, \dots, e'_{2n-1}$ for v' such that the normal connection

$$\omega'_{n+i-1, n+j-1} = 0 \quad (12)$$

Here and throughout this section, we shall agree on the index ranges

$$2 \leq i, j, k \leq n. \quad (13)$$

If we use condition (2) of the definition of a pseudo-spherical congruence, there is a local orthonormal frame field e_1, \dots, e_{2n-1} for M such that [9]

$$\left. \begin{aligned} e'_{n+i-1} &= -\sin \theta e_i + \cos \theta e_{n+i-1}, \\ e_1 &= \text{the unit direction of } \overrightarrow{PP'} \end{aligned} \right\} (14)$$

and e_1, \dots, e_n form an orthonormal frame for TM . Let

$$\left. \begin{aligned} e'_1 &= -e_1, \\ e'_i &= \cos \theta e_i + \sin \theta e_{n+i-1} \end{aligned} \right\} (15)$$

then e'_1, \dots, e'_n form an orthonormal frame for TM' . Since $\Gamma: v \rightarrow v'$ commutes with the normal connections, $\Gamma e_{n+i-1} = e'_{n+i-1}$, and $\omega'_{n+i-1, n+j-1} = 0$, we have

$$\omega_{n+i-1, n+j-1} = 0 \quad (16)$$

Suppose locally M is given by an immersion $X: U \rightarrow R^{2n-1}$, where U is an open subset of R^n , then M' is given by

$$X' = X + r e_1. \quad (17)$$

Taking the differential of (17) gives [9]

$$\left. \begin{aligned} dX' &= dX + r de_1, \\ &= \omega_1 e_1 + \sum_i \omega_i e_i + r \sum_i \omega_{1i} e_i + r \sum_i \omega_{1, n+i-1} e_{n+i-1} \\ &= \omega_1 e_1 + \sum_i (\omega_i + r \omega_{1i}) e_i + r \sum_i \omega_{1, n+i-1} e_{n+i-1}. \end{aligned} \right\} (18)$$

On the other hand, letting $\omega'_1, \dots, \omega'_n$ be the dual coframe of e'_1, \dots, e'_n , we have

$$\left. \begin{aligned} dX' &= \omega'_1 e'_1 + \sum_i \omega'_i e'_i, \quad \text{using(15)} \\ &= -\omega'_1 e_1 + \sum_i \omega'_i (-\cos \theta e_i + \sin \theta e_{n+i-1}). \end{aligned} \right\} (19)$$

Comparing coefficients of e_1, \dots, e_{2n-1} in (18) and (19), we get

$$\left. \begin{aligned} \omega'_1 &= -\omega_1, \\ \cos \theta \omega'_i &= \omega_i + r\omega_{1i}, \\ \sin \theta \omega'_i &= r\omega_{1,n+i-1}. \end{aligned} \right\} (20)$$

This gives

$$\omega_i + r\omega_{1i} = r \cot \theta \omega_{1,n+i-1}. \quad (21)$$

Using (12), (14) and (16), we have

$$\left. \begin{aligned} 0 &= \omega'_{n+i-1, n+j-1} \\ &= de'_{n+i-1} \cdot e'_{n+j-1} \\ &= d(-\sin \theta e_i + \cos \theta e_{n+i-1}) \cdot (-\sin \theta e_j + \cos \theta e_{n+j-1}) \\ &= \sin^2 \theta de_i \cdot e_j - \sin \theta \cos \theta (de_i \cdot e_{n+j-1} - de_j \cdot e_{n+i-1}) \\ &= \sin^2 \theta \omega_{ij} - \sin \theta \cos \theta (\omega_{i, n+j-1} - \omega_{j, n+i-1}). \end{aligned} \right\} (22)$$

Therefore we have

$$\omega_{ij} = \cot \theta (\omega_{i, n+j-1} - \omega_{j, n+i-1}). \quad (23)$$

In order to find the curvature, we compute the following 1-forms:

$$\left. \begin{aligned} \omega'_{1, n+k-1} &= de'_1 \cdot e'_{n+k-1}, \quad \text{using(14)and(15)} \\ &= -\sin \theta \omega_{1k} - \cos \theta \omega_{1, n+k-1}, \quad \text{using(21)} \\ &= -\frac{\sin \theta}{r} \omega_k, \\ \omega'_{i, n+k-1} &= de'_i \cdot e'_{n+k-1} \\ &= -\sin \theta \cos \theta \omega_{ik} + \cos^2 \theta \omega_{i, n+k-1} + \sin^2 \theta \omega_{k, n+i-1}, \quad \text{using(23)} \\ &= \omega_{k, n+i-1}. \end{aligned} \right\} (24)$$

Hence from equation (9) we have

$$\left. \begin{aligned} \Omega'_{1i} &= -\sum_k \omega'_{1, n+k-1} \wedge \omega'_{i, n+k-1}, \quad \text{using(24)} \\ &= \frac{\sin \theta}{r} \sum_k \omega_k \wedge \omega_{k, n+i-1}, \quad \text{using(1)} \\ &= -\frac{\sin \theta}{r} \omega_1 \wedge \omega_{1, n+i-1}, \quad \text{using(20)} \\ &= \frac{\sin^2 \theta}{r^2} \omega'_1 \wedge \omega'_i; \\ \Omega_{ij} &= -\sum_i \omega'_{i, n+k-1} \wedge \omega'_{j, n+k-1}, \quad \text{using(24)} \\ &= -\sum_i \omega_{k, n+i-1} \wedge \omega_{k, n+j-1}. \end{aligned} \right\} (25)$$

Since v is flat and $\omega_{n+i-1, n+j-1} = 0$, we have

$$\begin{aligned} 0 &= -d\omega_{n+i-1, n+j-1}, \quad \text{using(9)} \\ &= \omega_{1, n+i-1} \wedge \omega_{1, n+j-1} + \sum_k \omega_{k, n+i-1} \wedge \omega_{k, n+j-1}. \end{aligned}$$

So we have

$$\left. \begin{aligned} \Omega'_{1i} &= \omega_{1, n+i-1} \wedge \omega_{1, n+j-1}, \quad \text{using(9)} \\ &= \frac{\sin^2 \theta}{r^2} \omega'_i \omega'_j \end{aligned} \right\} (26)$$

Therefore M' has constant sectional curvature $-\sin^2\theta/r^2$. By symmetry, M also has constant sectional curvature $-\sin^2\theta/r^2$. [9]

Theorem 3.2

Suppose M is a local n -submanifold with constant negative sectional curvature $K = -\sin^2\theta/r^2$ in \mathbf{R}^{2n-1} , where $r > 0$ and θ are constants. Let v_1^0, \dots, v_n^0 , be an orthonormal base at P_0 consisting of principal curvature vectors, and $v_0 = \sum_{i=1}^n c_i v_i^0$ a unit vector with $c_i \neq 0$ for all $1 \leq i \leq n$; then there exists a local n -submanifold M' of \mathbf{R}^{2n-1} , and a pseudo-spherical congruence $l: M \rightarrow M'$ such that if $P'_0 = l(P_0)$, we have $\overline{P_0 P'_0} = r v_0$, and θ is the angle between the normal planes at P_0 and P'_0 . [9]

Theorem 3.3

Let $f_{\alpha i}, 1 \leq \alpha \leq 6, 1 \leq i \leq 3$, be differentiable functions of variables x, y and t such that [5,8]

$$\left. \begin{aligned} -f_{11,y} + f_{12,x} + \eta f_{52} + \xi f_{21} - \eta f_{22} - \xi f_{51} &= 0, \\ -f_{11,t} + f_{13,x} + \eta f_{53} + f_{43} f_{21} - \eta f_{23} - f_{51} f_{33} &= 0, \\ -f_{12,t} + f_{13,y} + \xi f_{53} + f_{43} f_{22} - f_{52} f_{33} - \xi f_{23} &= 0, \\ -f_{21,y} + f_{22,x} + \eta f_{62} + \eta f_{12} - \xi f_{61} - \xi f_{11} &= 0, \\ -f_{21,t} + f_{23,x} + \eta f_{63} + \eta f_{13} - f_{61} f_{33} - f_{11} f_{43} &= 0, \\ -f_{22,t} + f_{23,y} + \xi f_{63} + \xi f_{13} - f_{62} f_{33} - f_{12} f_{43} &= 0, \\ -f_{31,y} - f_{32,x} + f_{11} f_{52} - f_{12} f_{51} - f_{22} f_{61} + f_{21} f_{62} &= 0, \\ -f_{31,t} + f_{33,x} + f_{23} f_{61} + f_{13} f_{51} - f_{21} f_{63} - f_{11} f_{53} &= 0, \\ -f_{32,t} + f_{33,y} + f_{23} f_{62} + f_{13} f_{52} - f_{22} f_{63} - f_{12} f_{53} &= 0, \\ -f_{41,y} + f_{42,x} + f_{11} f_{22} - f_{12} f_{21} &= 0, \\ -f_{41,t} + f_{43,x} + f_{13} f_{21} - f_{11} f_{23} &= 0, \\ -f_{42,t} + f_{43,y} + f_{13} f_{22} - f_{12} f_{23} &= 0, \\ -f_{51,y} + f_{52,x} + \eta f_{12} - \xi f_{11} &= 0, \\ -f_{51,t} + f_{53,x} + \eta f_{13} - f_{11} f_{33} &= 0, \\ -f_{52,t} + f_{53,y} + \xi f_{13} - f_{12} f_{33} &= 0, \\ -f_{61,y} + f_{62,x} + \eta f_{22} - \xi f_{21} &= 0, \\ -f_{61,t} + f_{63,x} + \eta f_{23} - f_{21} f_{33} &= 0, \\ -f_{62,t} + f_{63,y} + \xi f_{23} - f_{22} f_{33} &= 0, \end{aligned} \right\} (27)$$

and

$$\begin{aligned} & f_{11} f_{22} - f_{12} f_{21} + \eta f_{12} - \xi f_{11} + \xi f_{51} - \eta f_{52} \\ & + \sin \phi [2\eta f_{12} - 2\xi f_{11} + f_{11} f_{52} - f_{12} f_{51} - \eta f_{52} - \xi f_{21} - f_{21} f_{52} + f_{21} f_{12} - f_{11} f_{22} + \xi f_{51} \\ & + \eta f_{22} + f_{22} f_{51}] \\ & + \cos \phi [f_{12} f_{51} - f_{11} f_{52} + 2\eta f_{12} - 2\xi f_{11} + \xi f_{21} + f_{21} f_{52} + f_{11} f_{22} - f_{21} f_{12} - \eta f_{52} \\ & - \eta f_{22} - f_{22} f_{51} + \xi f_{51}] = 0, \\ & f_{13} f_{21} - f_{11} f_{23} + \eta f_{13} - f_{11} f_{33} + f_{21} f_{33} + f_{33} f_{51} - \eta f_{53} - f_{21} f_{43} \\ & + \sin \phi [2\eta f_{13} - f_{11} f_{43} + f_{11} f_{53} - f_{13} f_{51} - \eta f_{43} - \eta f_{53} - f_{21} f_{43} - f_{21} f_{53} + f_{21} f_{13} + \eta f_{33} \\ & + f_{33} f_{51} - f_{11} f_{33} + \eta f_{23} + f_{23} f_{51} - f_{11} f_{23}] \\ & + \cos \phi [f_{13} f_{51} - f_{11} f_{53} + 2\eta f_{13} - f_{11} f_{43} + f_{21} f_{43} + f_{21} f_{53} - f_{21} f_{13} - \eta f_{43} - \eta f_{53} \\ & - \eta f_{23} - f_{23} f_{51} + f_{23} f_{11} + \eta f_{33} + f_{51} f_{33} - f_{11} f_{33}] = 0, \\ & f_{13} f_{22} - f_{12} f_{23} + \xi f_{13} - f_{12} f_{33} + f_{22} f_{33} + f_{52} f_{33} - f_{22} f_{43} - \xi f_{53} \\ & + \sin \phi [2\xi f_{13} - f_{12} f_{43} + f_{12} f_{53} - f_{13} f_{52} - \xi f_{43} - \xi f_{53} - f_{22} f_{43} - f_{22} f_{53} + f_{22} f_{13} + \xi f_{33} \\ & + f_{23} f_{52} - f_{23} f_{12} + f_{33} f_{52} - f_{33} f_{12} + \xi f_{23}] \\ & + \cos \phi [f_{13} f_{52} - f_{12} f_{53} + 2\xi f_{13} - f_{12} f_{43} + f_{22} f_{43} + f_{22} f_{53} - f_{22} f_{13} - \xi f_{43} - \xi f_{53} \\ & - \xi f_{23} - f_{23} f_{52} + f_{23} f_{12} + \xi f_{33} + f_{52} f_{33} - f_{12} f_{33}] = 0, \end{aligned}$$

With $f_{31} = \eta = f_{41}$; $f_{32} = \xi = f_{42}$

Then the following statements are valid.

1. The following system is completely integrable for ϕ ;

$$\left. \begin{aligned} 3\phi_x &= f_{41} + f_{51} + f_{61} - f_{11} + (f_{21} - f_{31}) \sin \phi + (f_{31} + f_{21}) \cos \phi \\ 3\phi_y &= f_{42} + f_{52} + f_{62} - f_{12} + (f_{22} - f_{32}) \sin \phi + (f_{32} + f_{22}) \cos \phi \\ 3\phi_t &= f_{43} + f_{53} + f_{63} - f_{13} + (f_{23} - f_{33}) \sin \phi + (f_{33} + f_{23}) \cos \phi \end{aligned} \right\} (28)$$
2. For any solution ϕ of(28) the 1- forms

$$\left. \begin{aligned} \sigma_1 &= f_{11} dx + f_{12} dy + f_{13} dt, \quad \text{and} \\ \sigma_2 &= (f_{21} \sin \phi + f_{31} \cos \phi) dx + (f_{22} \sin \phi + f_{32} \cos \phi) dy \\ &\quad + (f_{23} \sin \phi + f_{33} \cos \phi) dt \end{aligned} \right\} (29)$$

Are closed one – forms

3. If f_{ai} are analytic functions of parameters η and ξ at zero, then the solution $\phi(x, y, t, \eta, \xi)$ of (28) and the one-forms (29) are also analytic in η and ξ at zero.

Proof:

With respect to point 1. , it follows from the Frobenius theorem. and from (27) and (28) . straight forward computations show that (27) implies.

$$\phi_{xy} = \phi_{yx} \quad ; \quad \phi_{xt} = \phi_{tx} \quad ; \quad \phi_{yt} = \phi_{ty} .$$

point 2. , can be proved by showing that the systems (27) and (28) imply that exterior differentiation of the forms σ_1 and σ_2 in (29) is zero, which is the case.

In order to prove point 3. , we suppose that functions f_{ai} are analytic functions of parameters η and ξ . Each equation of (28) can be considered as an ordinary differential equation whose right – hand side is an analytic functions of (ϕ, η, ξ) , where the solutions of $\phi(x, y, t, \eta, \xi)$ of this equation exist as defined by point 1. It follows from the theory of ordinary differential equations, [11], on the dependence of solutions up on parameters, that $\phi(x, y, t, \eta, \xi)$ is an analytic functions of η and ξ , for η and ξ in an appropriate neighborhood of zero. This completes the proof of the theorem. [5,8]

IV. Derivation of Bäcklund transformations and conservation laws for evolution equations in higher dimensions

In this section, we extend the results obtained in [5], by introducing a new method to derive an infinite set of conservation laws for equations that describes a P.S.P., based on a geometrical property of these planes.

So, firstly. We consider M and M' as sub manifolds of R^{2n-1} of dim n , and $l: M \rightarrow M'$ be a pseudospherical geodesic congruence between M and M' , then there exist local orthonormal

forms [8,9] $e_1, e_2, \dots, e_n, e_{n+1}, \dots, e_{2n-1}$ and $e'_1, e'_2, \dots, e'_n, \dots, e'_{2n-1}$ for R^{2n-1} with e_1, e_2, \dots, e_n for M and e'_1, e'_2, \dots, e'_n for M' such that

$$\left. \begin{aligned} e'_1 &= \cos \theta e_1 + \sin \theta e_{n+1} , & 2 \leq i \leq n \\ e'_{n+i-1} &= -\sin \theta e_i + \cos \theta e_{n+i-1} , \end{aligned} \right\} (30)$$

Are verified, see [10], and $e'_1 = -e_1$, where e_1 at $P \in M$ is the unit vector tangent to the geodesic from P to $P' = l(P)$

In the special case, when $n = 2$, relations (30) become [1]

$$\left. \begin{aligned} e'_1 &= \cos \theta e_1 + \sin \theta e_2 \\ e'_2 &= -\sin \theta e_1 + \cos \theta e_2 \end{aligned} \right\} (31)$$

Where, it is considered that all the (n-1) angles are the same and equal to θ

In our case of evolution equations of three variables, M and M' 3- dimensional Riemannian sub manifolds of R^5 ; e'_1, e'_2, \dots, e'_5 and e_1, e_2, \dots, e_5 are two different orthonormal frames with e'_1, e'_2, e'_3 tangents to M' and e_1, e_2, e_3 tangents to M . While $\omega_1, \omega_2, \omega_3, \omega_{12}, \omega_{13}, \omega_{23}$ and $\omega'_1, \omega'_2, \omega'_3, \omega'_{12}, \omega'_{13}, \omega'_{23}$, are the dual coframes and connections forms on M and M' respectively. For $2 \leq i \leq 3(n = 3)$, one can write the following relations, with a pseudo spherical line congruence $l: M \rightarrow M'$,

$$\left. \begin{aligned} e'_1 &= -e_1 \\ e'_2 &= \cos \theta e_2 + \sin \theta e_3 \\ e'_3 &= -\sin \theta e_2 + \cos \theta e_3 \end{aligned} \right\} (32)$$

From the relations (32), we have

$$\left. \begin{aligned} \omega_1 &= -\omega'_1 \\ \omega_2 &= \cos \theta \omega'_2 + \sin \theta \omega'_3 \\ \omega_3 &= -\sin \theta \omega'_2 + \cos \theta \omega'_3 \end{aligned} \right\} (33)$$

And

$$\left. \begin{aligned} \omega'_1 &= -\omega_1 \\ \omega'_2 &= \cos \theta \omega_2 - \sin \theta \omega_3 \\ \omega'_3 &= \sin \theta \omega_2 + \cos \theta \omega_3 \end{aligned} \right\} (34)$$

Now, consider a differential equation (E) for $u(x, y, t)$ which describes a two parameters 3-dim. P.S.P. with associated 1 – forms.

$$\left. \begin{aligned} \omega_1 &= f_{11}dx + f_{12}dy + f_{13}dt; \\ \omega_2 &= f_{21}dx + f_{22}dy + f_{23}dt; \\ \omega_3 &= \eta dx + \xi dy + f_{33}dt; \\ \omega_4 &= \omega_{12} = \eta dx + \xi dy + f_{43}dt; \\ \omega_5 &= \omega_{13} = f_{51}dx + f_{52}dy + f_{53}dt; \\ \omega_6 &= \omega_{23} = f_{61}dx + f_{62}dy + f_{63}dt. \end{aligned} \right\} (35)$$

Where f_{ai} , are functions of $u(x, y, t)$ and its derivatives (observe that we are denoting ω_4, ω_5 and ω_6 for ω_{12}, ω_{13} and ω_{23} respectively, which are the classical notation for the connection forms). We have the following[8]

Proposition 4.1

Let E be a differential equation which describes a two parameters 3- dim P.S.P. with associated 1- forms (35) . Then, for each solution u of E , the system of equations for $\phi(x, y, t)$. [8]

$$\left. \begin{aligned} \omega_{12} - d\phi + \omega'_3 &= 0, \\ \omega_{13} - d\phi + \omega'_2 &= 0, \\ \omega_{23} - d\phi + \omega'_1 &= 0 \end{aligned} \right\} (36)$$

Are completely integrable. Moreover, for each solution u of E , and corresponding solution ϕ of (36). the forms

$$\left. \begin{aligned} \sigma_1 &= f_{11}dx + f_{12}dy + f_{13}dt, \\ \sigma_2 &= (f_{21} \sin \phi + f_{31} \cos \phi)dx + (f_{22} \sin \phi + f_{32} \cos \phi)dy + (f_{23} \sin \phi + f_{33} \cos \phi)dt \end{aligned} \right\}$$

Are closed forms.

Proof

It follows from (33) that u is a solution of E iff. $\omega_{12} - d\phi + \omega'_3 = 0$ i.e.,

$$\left. \begin{aligned} \omega_{12} - d\phi + \sin \phi \omega_2 + \cos \phi \omega_3 &= 0, \\ \omega_{13} - d\phi + \omega'_2 &= 0 \end{aligned} \right\} (37)$$

$$\left. \begin{aligned} \omega_{13} - d\phi + \cos \phi \omega_2 + \sin \phi \omega_3 &= 0, \\ \omega_{23} - d\phi + \omega'_1 &= 0 \end{aligned} \right\} (38)$$

$$\text{and } \omega_{12} - d\phi + \omega_1 = 0 (39)$$

Are completely integrable for ϕ . In this case:

$$\omega_1 ; \text{ and } \sin \phi \omega_2 + \cos \phi \omega_3 (40)$$

Are closed forms. Hence, inserting (35) into (37) we obtain equations.

$$(\eta - \phi_x + f_{21} \sin \phi + \eta \cos \phi)dx + (\xi - \phi_y + f_{22} \sin \phi + \xi \cos \phi)dy + (f_{43} - \phi_t + f_{23} \sin \phi + f_{23} \cos \phi)dt = 0$$

$$3\phi_x = f_{41} + f_{51} + f_{61} - f_{11} + (f_{21} - f_{31}) \sin \phi + (f_{31} + f_{21}) \cos \phi, (41)$$

Are closed forms. Hence, inserting (35) into (38) we obtain equations.

$$(f_{51} - \phi_x + f_{21} \cos \phi - \eta \sin \phi)dx + (f_{52} - \phi_y + f_{22} \cos \phi - \xi \sin \phi)dy + (f_{53} - \phi_t + f_{23} \cos \phi - f_{23} \sin \phi)dt = 0$$

$$3\phi_y = f_{42} + f_{52} + f_{62} - f_{12} + (f_{22} - f_{32}) \sin \phi + (f_{32} + f_{22}) \cos \phi, (42)$$

Are closed forms. Hence, inserting (35) into (39) we obtain equations.

$$(f_{61} - \phi_x - f_{11})dx + (f_{62} - \phi_y - f_{12})dy + (f_{63} - \phi_t - f_{13})dt = 0$$

$$3\phi_t = f_{43} + f_{53} + f_{63} - f_{13} + (f_{23} - f_{33}) \sin \phi + (f_{33} + f_{23}) \cos \phi (43)$$

Whose integrability condition is E . also, inserting (35) into (40), One can obtain the closed forms (29).

Now, we note that whenever E does not involve the parameters η, ξ the closed forms (29) may provide an infinite number of conservation laws. [8]

Also, under certain conditions, equations (28) may provide Bäcklund transformations for E and as we know from theorem (3.3), the conditions

$$\phi_{xy} = \phi_{yx} \quad ; \quad \phi_{xt} = \phi_{tx} \quad ; \quad \phi_{yt} = \phi_{ty} \quad .$$

are valid as the complete integrability condition for (28). As conservation laws are common features of mathematical physics, where they describe the conservation of fundamental physical quantities, it is worth studying them in this geometric study. Before giving the method for deriving conservation laws for evolution equations that describes 2 – parameters 3-dim. P.S.P., we consider the following:

Definition 4.1

We suppose a system of the form

$$u_t = S[u] (44)$$

In the system, when a functional $J[(x, y, t)]$ satisfies.

$$dJ[U(x, y, t)]/dt = 0, (45)$$

The functional is said to be an integral of equation (44). and

$$\frac{\partial}{\partial t} T[u(x, y, t)] + \frac{\partial}{\partial x} Q[u(x, y, t)] + \frac{\partial}{\partial y} R[u(x, y, t)] = 0 \quad (46)$$

Where usually each of $T[u(x, y, t)]$, $Q[u(x, y, t)]$ and $R[u(x, y, t)]$ do not involve derivatives with respect to t , is called a conservation law. In particular, if we are to apply this idea to an evolution equations for $u(x, y, t)$, then T, Q and R may depend upon $x, y, t, u, u_x, u_y, u_{xx}, u_{yy}, \dots$, but not on u_t . [8]

If we assume the function $u(x, y, t)$ and its derivatives with respect to x and y go to zero sufficiently fast as $|x| \rightarrow \infty, |y| \rightarrow \infty$, i. e., if T, Q_x and R_y are integrable on $(-\infty, \infty)$, so that $Q \rightarrow \text{constant}$ as $|x| \rightarrow \infty, R \rightarrow \text{constant}$ as $|y| \rightarrow \infty$, then equation (46) can be integrated to yield.

$$\frac{d}{dt} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} T dx dy = 0, \text{ i. e. } \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} T[u(x, y, t)] dx dy = J(u) = \text{Constant}$$

The method to derive conservation laws for evolution equations that describe, spherical surfaces (P.S.S.) is introduced by Cavalcante and Tenenblatin [5], for the case of two independent variables.

In this work, we will give a method with argument analogue to that considered in [5,8] to derive conservation laws for evolution equations that describe 2 – parameters 3-dim. P.S.P. This integrated method is based on geometrical properties of these planes.

Here, we suppose the functions f_{ai} to be analytic in each of η and ξ separately, and describe the solutions ϕ of (28) as a power series of η and ξ . In addition, from relations (29) we obtain a sequence of closed one – forms. So, we suppose

$$f_{ai}(x, y, t, \eta, \xi) = \sum_{r=0}^{\infty} h_{ai}^r(x, t) \eta^r + g_{ai}^r(y, t) \xi^r \quad (47)$$

And the solution ϕ of (28) may have the form

$$\phi(x, y, t, \eta, \xi) = \sum_{i=0}^{\infty} (\phi_i(x, t) \eta^i + \psi_i(y, t) \xi^i) \quad (48)$$

For fixed y, t , we consider functions of η and ξ respectively as follows:

$$C(\eta, \xi) = \cos \phi = \cos \left[\sum_{i=0}^{\infty} (\phi_i \eta^i + \psi_i \xi^i) \right] \quad (49)$$

$$S(\eta, \xi) = \sin \phi = \sin \left[\sum_{i=0}^{\infty} (\phi_i \eta^i + \psi_i \xi^i) \right] \quad (50)$$

From relations (49) and (50), we have

$$\left. \begin{aligned} C(0,0) &= \cos(\phi_0 + \psi_0); \\ S(0,0) &= \sin(\phi_0 + \psi_0) \end{aligned} \right\} \frac{d^r C}{d\eta^r}(0, \xi) = -(r-1)! \sum_{\alpha=0}^{r-1} \frac{r-\alpha}{\alpha!} \frac{d^\alpha C}{d\eta^\alpha}(0, \xi) \phi_{r-\alpha} \quad (51)$$

$$\frac{d^r}{d\eta^r}(0, \xi) = (r-1)! \sum_{\alpha=0}^{r-1} \frac{r-\alpha}{\alpha!} \frac{d^\alpha C}{d\eta^\alpha}(0, \xi) \phi_{r-\alpha}, \text{ for } r \geq 1$$

$$\frac{d^r C}{d\xi^r}(\eta, 0) = -(r-1)! \sum_{\alpha=0}^{r-1} \frac{r-\alpha}{\alpha!} \frac{d^\alpha S}{d\xi^\alpha}(\eta, 0) \psi_{r-\alpha} \quad (52)$$

$$\frac{d^r S}{d\xi^r}(\eta, 0) = (r-1)! \sum_{\alpha=0}^{r-1} \frac{r-\alpha}{\alpha!} \frac{d^\alpha C}{d\xi^\alpha}(\eta, 0) \psi_{r-\alpha}, \text{ for } r \geq 1$$

Finally, we define the following functions of x, y, t :

$$G_r^{\alpha i} = (h_{2r}^\alpha - h_{3r}^\alpha) \frac{d^{i-\alpha} C}{d\eta^{i-\alpha}}(0, \xi) - (h_{2r}^\alpha + h_{3r}^\alpha) \frac{d^{i-\alpha} S}{d\eta^{i-\alpha}}(0, \xi),$$

$$L_r^{\alpha i} = (h_{2r}^\alpha - h_{3r}^\alpha) \frac{d^{i-\alpha} S}{d\eta^{i-\alpha}}(0, \xi) + (h_{2r}^\alpha + h_{3r}^\alpha) \frac{d^{i-\alpha} C}{d\eta^{i-\alpha}}(0, \xi),$$

$$F_{3r} = (h_{4r}^3 + h_{5r}^3 + h_{6r}^3 - h_{1r}^3) + L_r^{33}, \quad r = 1, 3$$

$$F_{Pr} = (h_{4r}^P + h_{5r}^P + h_{6r}^P - h_{1r}^P) + \sum_{m=1}^{P-1} \frac{P-m}{m!} G_r^{0m} \phi_{P-m} + \sum_{m=1}^P \frac{1}{(P-m)!} L_r^{mP}, \quad r = 1, 3$$

$$'G_r^{\alpha i} = (g_{2r}^\alpha - g_{3r}^\alpha) \frac{d^{i-\alpha} C}{d\xi^{i-\alpha}}(\eta, 0) - (g_{2r}^\alpha + g_{3r}^\alpha) \frac{d^{i-\alpha} S}{d\xi^{i-\alpha}}(\eta, 0),$$

$${}'L_r^{\alpha i} = (g_{2r}^\alpha - g_{3r}^\alpha) \frac{d^{i-\alpha} S}{d\xi^{i-\alpha}}(\eta, 0) + (g_{2r}^\alpha + g_{3r}^\alpha) \frac{d^{i-\alpha} C}{d\xi^{i-\alpha}}(\eta, 0),$$

$$F'_{3r} = (g_{4r}^3 + g_{5r}^3 + g_{6r}^3 - g_{1r}^3) + L_r^{33}$$

$$F'_{qr} = (g_{4r}^q + g_{5r}^q + g_{6r}^q - g_{1r}^q) + \sum_{s=1}^{q-1} \frac{q-s}{s!} G_r^{0s} \psi_{q-s} + \sum_{s=1}^q \frac{1}{(q-s)!} L_r^{sq}, \quad r = 2, 3$$

Where each of α, i, p, q is non-negative integer such that $i \geq \alpha, q \geq 4$, and $r = 1, 2, 3$.

It is easy to see that the functions $G_r^{\alpha i}$ and $L_r^{\alpha i}$ depend above depend on $\phi_0, \phi_1, \dots, \phi_{i-\alpha}$. Whereas the functions $'G_r^{\alpha i}$ and $'L_r^{\alpha i}$ depend on $\psi_0, \psi_1, \dots, \psi_{i-\alpha}$

Also, the functions (F_{2r}, F_{3r}) and (F_{pr}) depend on ϕ_0 and $\phi_1, \phi_2, \dots, \phi_{p-1}$, respectively, but the functions (F'_{2r}, F'_{3r}) and (F'_{qr}) depend on ψ_0 and $\psi_1, \psi_2, \dots, \psi_{q-1}$, respectively.

Under the above notation, we obtain the following corollary.

Corollary 4.1

Suppose $f_{\alpha i}(x, y, t, \eta, \xi)$, $1 \leq \alpha \leq 6$, $1 \leq i \leq 3$, be differentiable functions of x, y, t , analytic at $\eta = 0, \xi = 0$ that satisfy (27). Then, in view of the above notation, the following statements hold.

(i) The solution ϕ of (28) is analytic at $\eta = 0, \xi = 0$; ϕ_0 and ψ_0 are determined by.

And, for $1 \geq i, \phi_i$ and ψ_i are recursively determined by the system

(ii) For any such solution ϕ , equation (48) and any integer $i \geq 0$

Are closed one-forms.

The proof of the corollary follows with somehow straight forward calculations from equations (48) \rightarrow (54) and equations (28) with the introduced notations.

Now, if we consider a non-linear evolution equation for $u(x, y, t)$ which describes a 3-dim.

P.S.P., then there exist functions $f_{\alpha i}, 1 \leq \alpha \leq 6, 1 \leq i \leq 3$, depending on $u(x, y, t)$ and its derivatives, such that, for any solution u of the evolution equation, $f_{\alpha i}$ satisfy (27). So, it follows from theorem (3) that equations (28) are completely integrable for . If we consider $f_{\alpha i}$ to be analytic functions of parameters η, ξ then we can find that the solutions ϕ of (28) and the 1-forms given by (29), are analytic in η, ξ where their coefficients ϕ_i, ψ_i and β^i , as functions of u , are determined by (55) \rightarrow (59). The closed 1-forms β^i provide a sequence of conservation laws for the evolution equation, with equations given by.

$$Q_i = \sum_{\alpha=0}^i \frac{1}{(i-\alpha)!} (G_3^{\alpha i} + {}'G_3^{\alpha i} + {}'G_2^{\alpha i}),$$

$$T_i = \sum_{\alpha=0}^i \frac{1}{(i-\alpha)!} (G_1^{\alpha i} + G_2^{\alpha i}),$$

with

$$Q_{i,x} + R_{i,x} + T_{i,1} = 0, \quad i \geq 0$$

For Bäcklund transformations of the equation E which describes P.S.P., we remark that the angle ϕ of a pseudospherical line congruence is determined by the system of equations (28). If we suppose that (28) is equivalent to a system of the form:

Then given a solution u of E the system (28) is integrable and ϕ is a solution of (62), then u defined by the equation (61) will be a solution for E . However, it still needs more work to be done.

V. Conclusion

In this paper, we generalized Bäcklund transformations and conservation laws for evolution equations in higher dimensions.

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References

- [1]. Chern. S.S and Tenenblat. K, " pseudospherical surfaces and soliton equations" Stud. Appl. Mat. 74, 55-83 (1986).
- [2]. El-Sabbagh. M and K.R.Abdo, "Pseudospherical surfaces and evolution equations in higher dimensions "IJSET. Math. vol (4) Issue No.3, 165-171 (2015).
- [3]. El-Sabbagh. M and K.R.Abdo, "Pseudospherical planes and evolution equations in higher dimensions II" IOSR-JM. Math. vol(11) Issue No.2, Ver. I, 102-111(2015).
- [4]. El-Sabbagh. M and K.R.Abdo, "Pseudospherical 3- Planes In \mathbb{R}^5 and Evolution Equations Of Type $u_{tt} = \psi \left(u, u_x, \dots, \frac{\partial^k u}{\partial x^k}, u_y, \dots, \frac{\partial^k u}{\partial y^k}, u_t \right)$ " IOSR-JM. Math. Vol (11) Issue No.2, Ver. PP 28-38(2015).
- [5]. J. A. Cavalcante and K. Tenenblat, "Conservation laws for nonlinear evolution equations" Stud. Appl. Math. 74, 1. 1044-1049(1988).
- [6]. Rabelo, M. L.; "On evolution equations which describe pseudo-spherical surfaces", Stud. Appl. Math.81, 221-248 (1989).
- [7]. Rabelo M.L and Tenenblat K., " On equations of the type $u_{xt} = F(u, u_x)$ which describe Pseudospherical surfaces", J. Mat, phys. 31(6), 1400-1407 (1990).
- [8]. El-Sabbagh. M, " Bäcklunds Transformations and Conservation Laws For evolution equations in higher dimensions " EL-Minia Science Bulletin, vol. 6, No.6, (1993).
- [9]. Tenenblat. K. and Treng .C, "Bäcklunds theorems for n-dimensional Submanifolds of \mathbb{R}^{2n-1} ", Ann. of Mat. vol. 111, 477-490 (1980).
- [10]. Tenenblat. K.;" Backlund's Theorem For Submanifolds Of Space forms And A Generalized Wave Equation" Bol. Soc. Bras. Mat., Vol. 16 N2, 67-92(1985).
- [11]. E. Coddington and N. Levinson, " Theory of ordinary differential equations", McGraw Hill, New York (1995).
- [12]. P. Bennett , J. Math. vol 31, No. 12 (1990) 2872-2916.