

A Study of Longitudinal Surface Roughness Effect on the performance of infinitely short Journal Bearing

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Abstract:

This paper attempts to study the performance of an infinitely short rough journal bearing. The study has been extracted using the Christenson and Tonder model. The corresponding Reynolds equation has been used to calculate the pressure distribution, which is important for constructing the load-carrying capacity equation. A graphical representation is made to display the calculated value of dimensionless pressure and dimensionless load-carrying capacity. It can be underlined that the standard deviation significantly affects the performance characteristics.

Keywords: Journal Bearing, Longitudinal surface roughness, Reynolds equation, Pressure, Load carrying capacity.

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I. Introduction

In recent decades, hydrodynamic journal bearings have received great attention from practical and analytical point of view. A rapid growth of sleeve bearing technology is mainly due to its wide range of engineering applications such as precision machine tools, nuclear reactors, high-speed aircraft and textile spindles. The development of thin-film lubrication provided a theoretical basis for revealing practical engineering.

Kuzma¹, Lin², Nada and Osman³, Motazeri⁴, Patel et al.⁵ and Lin et al.⁶ have studied on hydrodynamic journal bearings assuming a surface to be smooth. Several works have found on the performance of the journal bearings, such as the works in short bearings by Lin⁷, Naduvinamani et al.⁸ Patel et al.⁹ and Lin et al.¹⁰; long bearings by Lin¹⁰, Shah and Bhatt¹¹ and Naduvinamani et al.¹²; infinite bearings by Gururajan and Prakash¹³; narrow bearing by Gururajan and Prakash¹³ and finite bearings by Lin¹⁴, Elsharkawy and Guedouar¹⁵ and Elsharkawy¹⁶.

In almost all types of bearing systems, a very important role is played by the roughness of the bearing surface and its roughness pattern. Some investigators (namely Michell¹⁷, Pinkus and Sternlicht¹⁸, Burton¹⁹, Davies²⁰, Tzeng and Saibel²¹ have observed that the performance of the bearing and its life depend on the roughness of the surface, friction, wear and lubricants.

After that, many investigators viz. Andharia et al.^{22,23,24}, Shah and Bhat²⁷, Deheri et al.^{25,26,28}, Patel et al.^{29,30}, Shukla and Dehari³¹, Naduvinamani et al.³⁴, Andharia and Patel^{32,33}, Panchal et al.^{35,36} have studied and established that the performance of the bearing depends on its roughness and they have also given essential results in terms of the load-carrying capacity of the bearing for various parameters such as the mean, standard deviation and asymmetry of the bearing with and without the use of a magnetic lubricant. In these investigations, Christensen and Tender's^{37,38,39}, Christensen⁴⁰ role could not be underestimated, who produced a stochastically averaged Reynolds type equation by applying a random averaging process on a Reynolds type equation that gives the mean pressure more readily than local pressure.

Recently, more and more consideration has been given to study the effect of longitudinal surface roughness by Andharia and Deheri^{41,42,43}, Sharma and Pandey⁴⁴, Andharia and Patel⁴⁵, Patel and Deheri⁴⁶, Lin⁴⁷, Shimpi and Deheri⁴⁸, Andharia et al.⁴⁹, Andharia and Pandya^{50,51,52}, Adeshara et al.⁵³ and Patel et al.⁵⁴. All of these observations above clearly show that from a bearing longevity point of view, the effect of the standard deviation remains very crucial.

In this paper, an averaged Reynolds-type stochastic equation for infinitely short journal bearing, and the rough longitudinal bearing surface is solved numerically, and graphical results are produced accordingly.

Indeed, it is revealed that the load-carrying capability of the infinitely short journal bearing can be significantly improved.

II. Analysis:

The pressure distribution along circumference is shown in figure 1. The assumptions of general hydrodynamic lubrication theory are taken into consideration in the development of the analysis.

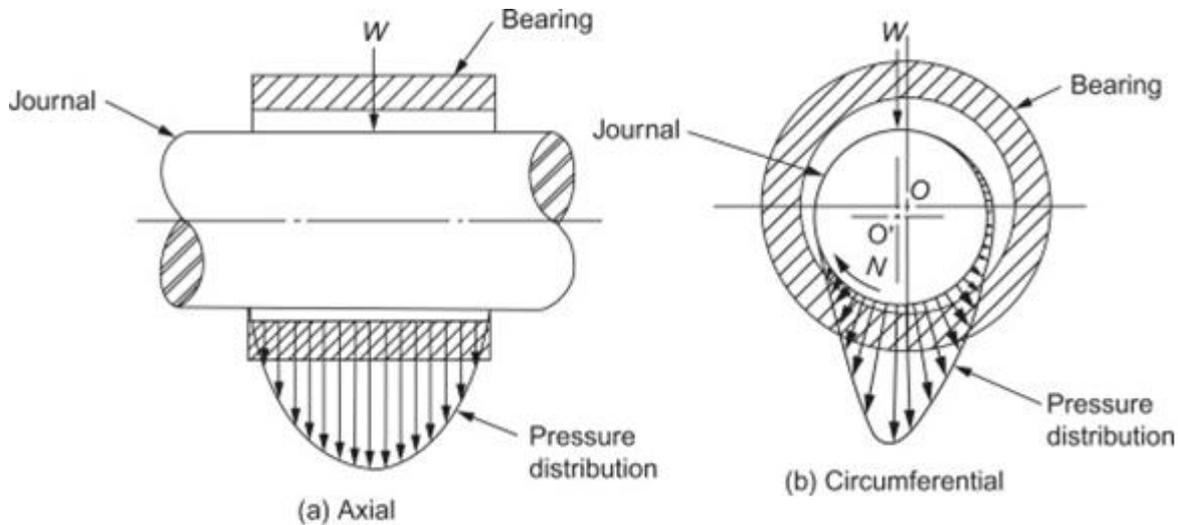


Figure: 1 Physical configuration of a journal bearing

The pressure governing equation is the Reynolds equation in two dimensions for an incompressible fluid, given by

$$\frac{\partial}{\partial x} \left(h^3 \frac{\partial p}{\partial x} \right) + \frac{\partial}{\partial z} \left(h^3 \frac{\partial p}{\partial z} \right) = 6 \eta U \frac{dh}{dx} \quad (1)$$

The short bearing approximation assumes a sufficiently short bearing in the axial direction and for infinitely short journal bearing

$$\frac{\partial p}{\partial x} = 0. \quad (2)$$

Substitute the value of (2) in equation (1), The Reynolds equation for a short journal bearing is obtained as

$$\frac{\partial}{\partial z} \left(h^3 \frac{\partial p}{\partial z} \right) = 6 \eta U \frac{dh}{dx} \quad (3)$$

The bearing surfaces are assumed to be longitudinally rough. Following Christensen and Tonder (1969, 1971 and 1972) the expression for film thickness is considered as

$$h(x) = \bar{h}(x) + \delta_s$$

where h is the mean film thickness, δ_s is the random deviation from the mean film thickness governed by probability density function $f(\delta_s)$, $-c \leq \delta_s \leq c$ [Andharia et al. (1997)]. Here c being the maximum deviation in δ_s .

According to the stochastic averaging method for longitudinal roughness discussed in Christensen and Tonder (1969, 1971, 1972), Andharia et al. (1997) Eq.(3) is transformed to

$$\frac{\partial}{\partial z} \left(\frac{1}{E(h^{-3})} \frac{\partial p}{\partial z} \right) = 6 \eta U \frac{d}{dx} \frac{1}{E(h^{-1})} \quad (4)$$

where, $E(h^{-3}) = \bar{h}^{-3}(1 + 6\bar{h}^{-2}\sigma^2)$
 $E(h^{-1}) = \bar{h}^{-1}(1 + \bar{h}^{-2}\sigma^2).$

Hence,

$$\frac{\partial}{\partial z} \left(\frac{1}{\bar{h}^{-3}(1+6\bar{h}^{-2}\sigma^2)} \frac{\partial p}{\partial z} \right) = 6 \eta U \frac{d}{dx} \left(\frac{1}{\bar{h}^{-1}(1+\bar{h}^{-2}\sigma^2)} \right). \quad (5)$$

Let $x = r\theta$ and hence: $dx = r d\theta$ and the film thickness h is described by the following approximation, which is used in almost all analyses: $\bar{h} = c + e \cos\theta = c(1 + \varepsilon \cos\theta)$ that gives $\frac{d\bar{h}}{d\theta} = -\varepsilon \sin\theta$, where (eccentricity) $E = \frac{e}{c}$. Fluid film thickness \bar{h} is not a function of z ; hence equation (5) becomes

$$\frac{d^2 p}{dz^2} = 6 \eta U \left\{ \bar{h}^{-3}(1 + 6\bar{h}^{-2}\sigma^2) \right\} \frac{d}{rdx} \left(\bar{h}^{-1}(1 + \bar{h}^{-2}\sigma^2) \right)^{-1} \quad (6)$$

Integrating the above equation twice with respect to z , Eq. (6) takes the form

$$p = 3 \eta U \left\{ \bar{h}^{-3}(1 + 6\bar{h}^{-2}\sigma^2) \right\} \frac{d}{rdx} \left(\bar{h}^{-1}(1 + \bar{h}^{-2}\sigma^2) \right)^{-1} z^2 + Cz + D \quad (7)$$

where C and D are arbitrary constant.

Using the boundary conditions at the point $z = \pm \frac{L}{2}$, $p = 0$ and $z = 0$, $\frac{dp}{dz} = 0$ in Eq. (7) and simplifying, one get the expression for pressure distribution in the bearing system as

$$p = \frac{3 \eta U L^2 (-E \sin \theta)}{c^2 r (1+E \cos \theta)^3} \left\{ \frac{\left(1+3\left(\frac{\sigma}{c}\right)^2 \frac{1}{(1+E \cos \theta)^2}\right)}{\left(1+\left(\frac{\sigma}{c}\right)^2 \frac{1}{(1+E \cos \theta)^2}\right)^2} \right\} \left(1+6\left(\frac{\sigma}{c}\right)^2 \frac{1}{(1+E \cos \theta)^2}\right) \left[\frac{z^2}{L^2} - 0.25\right] \quad (8)$$

Introducing of dimensionless quantities

$$\sigma^* = \frac{\sigma}{c}, z^* = \frac{z}{L}, P^* = -\frac{C^2 r p}{\eta U L^2},$$

in Eq. (8) it leads to the expression of the pressure distribution in the dimensionless form as

$$P^* = \frac{3 (E \sin \theta)}{(1+E \cos \theta)^3} \left\{ \frac{\left(1+\frac{3(\sigma^*)^2}{(1+E \cos \theta)^2}\right)\left(1+\frac{6(\sigma^*)^2}{(1+E \cos \theta)^2}\right)}{\left(1+\frac{(\sigma^*)^2}{(1+E \cos \theta)^2}\right)^2} \right\} [(z^*)^2 - 0.25]. \quad (9)$$

The component of the total force in the direction of the line of centres is described by,

$$W_N^* = \int_0^\pi \int_{-\frac{1}{2}}^{\frac{1}{2}} P^* \cos \theta \, d\theta \, dz^*.$$

The component of the total force in the direction of the line of centres is given by,

$$W_T^* = \int_0^\pi \int_{-\frac{1}{2}}^{\frac{1}{2}} P^* \sin \theta \, d\theta \, dz^*.$$

Load – bearing capacity of the bearing is given by

$$W^* = \sqrt{W_N^{*2} + W_T^{*2}}. \quad (10)$$

Resulting load-bearing capacity is computed by using Simpson's 1/3 rule.

III. Results and Discussion

The theoretical aspects of the results calculated numerically for various bearing characteristics are explained by figures. In order to study a quantitative behavior of the inconsistency variations of the incompressible lubricant, it is first necessary to calculate the pressure and the load capacity. This is achieved by solving the simultaneous Eq. (5) for the dimensionless pressure P^* and the Eq. (9) for the dimensionless load capacity W^* by the 1/3 method of Simpson. The variation of the pressure distribution (P^*) is shown in Fig. – 2 to Fig. – 6 and the load capacity (W^*) is shown in Fig. – 7 to Fig. – 8.

The pressure distribution P^* versus θ for various values of Z^* , E and σ^* has been depicted in Fig. – 2 to Fig. – 4. The pressure distribution (P^*) decreases continuously as θ increases.

The pressure distribution P^* versus Z^* for various values of E and σ^* has been depicted in fig-5 and fig-6. The pressure distribution (P^*) decreases continuously as Z^* increases.

The load capacity W^* versus σ^* for the variable E has been shown in Fig. – 7. The load capacity is continuously increasing to increase the value of E .

The load capacity W^* versus E for the variable σ^* has been shown in Fig. – 8. The load capacity increases sharply as the value of σ^* increases.

Here, it is depicted that the increasing value of the standard deviation (σ^*) and E increases the load capacity.

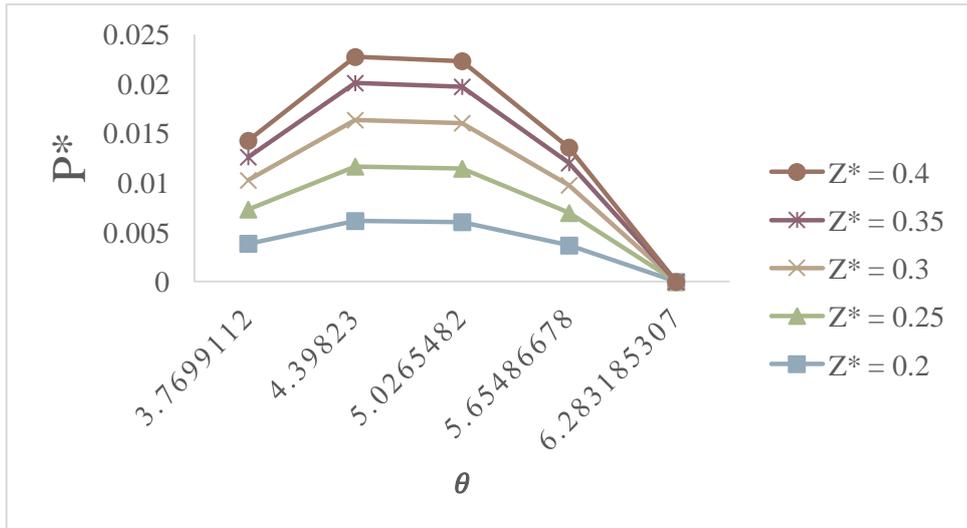


Figure – 2: Effect on pressure distribution with respect to θ for variable Z^* .

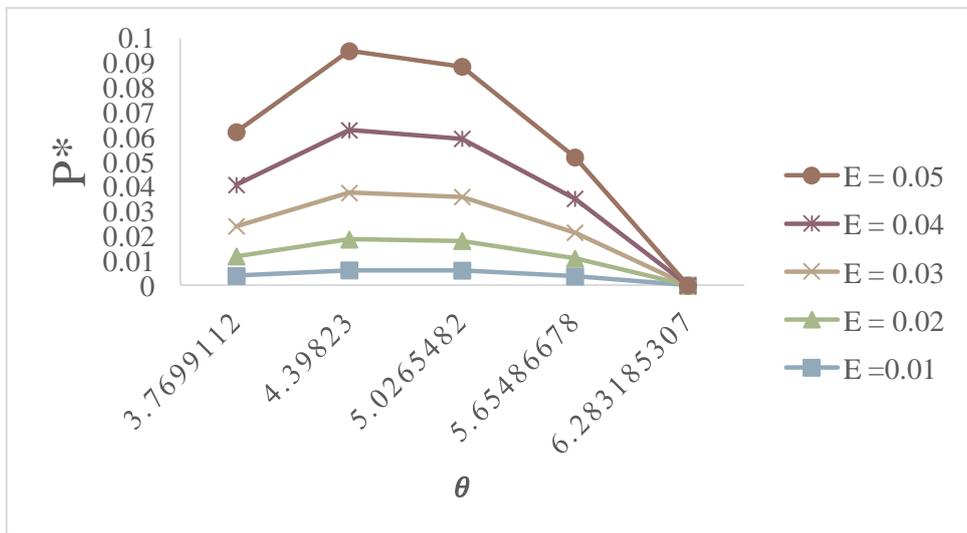


Figure – 3: Effect on pressure distribution with respect to θ for variable E .

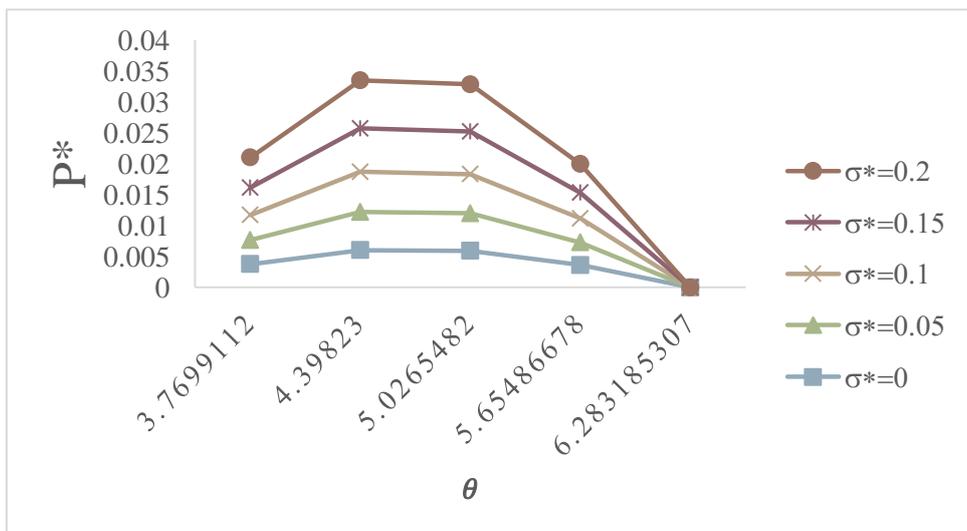


Figure – 4: Effect on pressure distribution with respect to θ for variable σ^* .

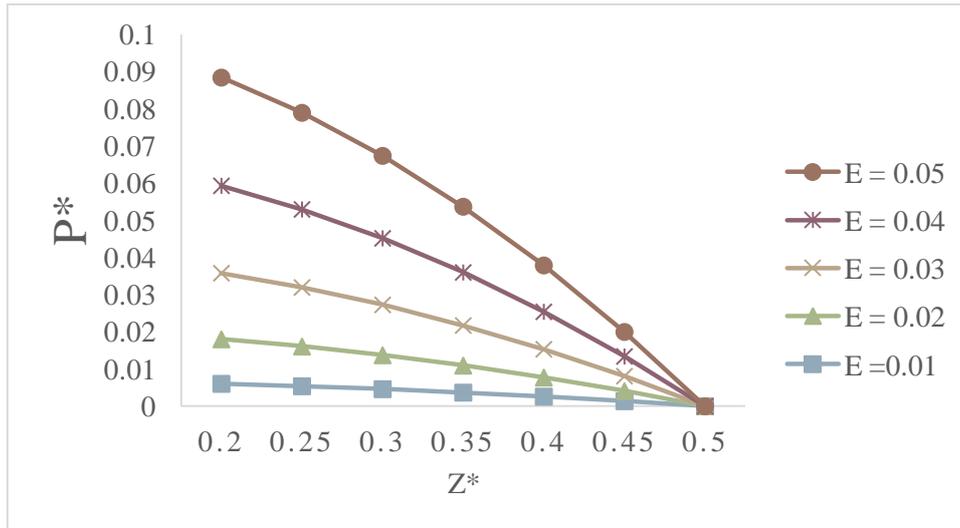


Figure – 5: Effect on pressure distribution with respect to Z^* for variable E .

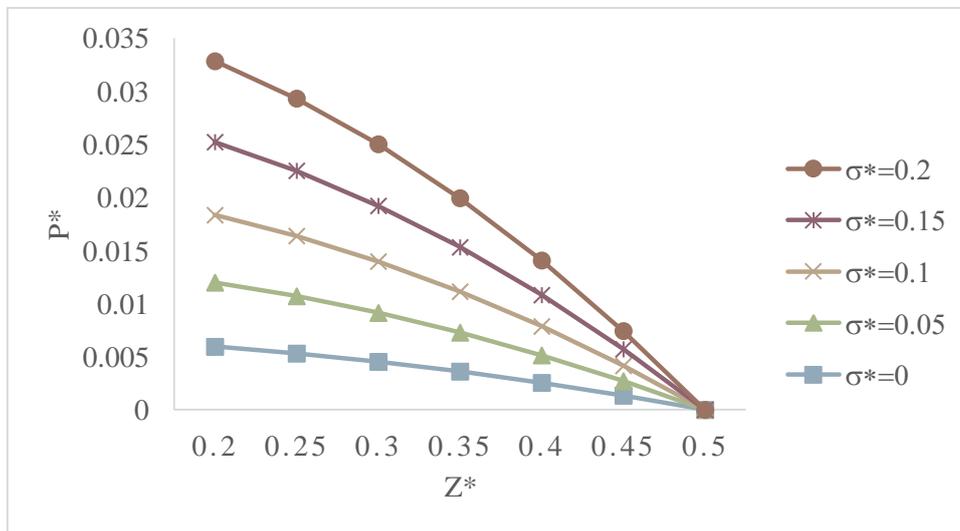


Figure – 6: Effect on pressure distribution with respect to Z^* for variable σ^* .

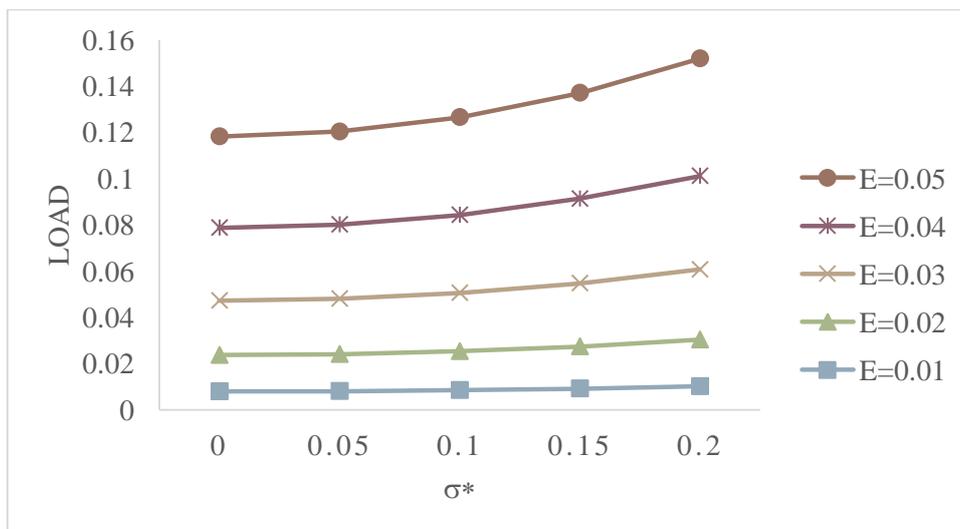


Figure – 7: Effect on load-carrying capacity with respect to σ^* for variable E .

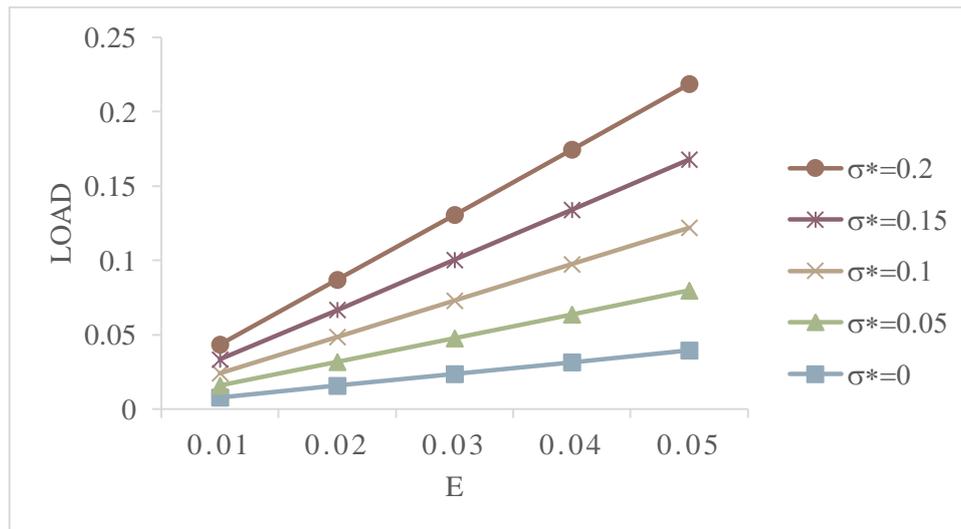


Figure – 8: Effect on load-carrying capacity with respect to E for variable σ^* .

IV. Conclusions:

This study emphasizes that roughness aspects must be evaluated when designing this type of bearing system which is important from bearing life. It is often observed that the standard deviation and eccentricity greatly improve load-carrying capacity.

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